
HydDown

User guide and technical reference

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1 Introduction

HydDown is an open source python tool for calculation of Hydrogen (or other pure gas phase species) vessel/container depressurization and filling. The HydDown logo shown in Fig. 1.1 visualizes the key parameters and transport phenomena during gas vessel filling or discharging. The thermodynamic state inside the vessel changes over time as seen from immediately observable variables temperature (T) and pressure (P). This is caused by change in fluid inventory (density) due to flow of gas either in- or out of the vessel. Further, heat is transferred from or to the surroundings via convective heat transfer on the in- and outside of the vessel - with heat being conducted through the vessel wall.

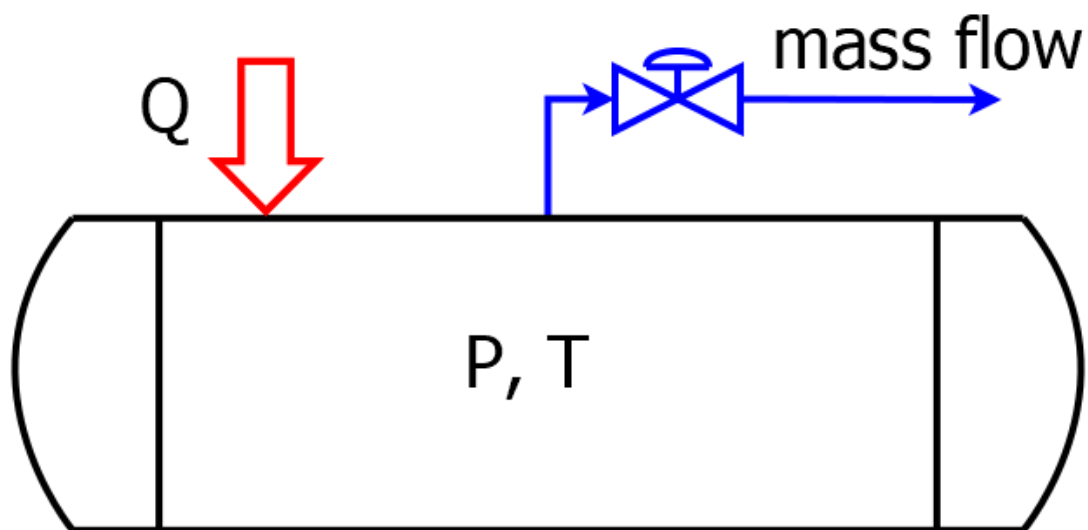


Figure 1.1: HydDown logo

Run the code as simple as:

```
1 python main.py input.yml
```

where main.py is the main script and input.yml is the input file in Yaml syntax.

1.1 Background

This is a small spare time project for calculation of vessel filling and depressurisation behaviour. This is mainly to demonstrate, that although perceived as a very tedious/difficult task to write your own code for such an apparent complex problem, actually a fairly limited amount of code is necessary if you have a good thermodynamic backend.

A few choices are made to keep things simple to begin with:

- Coolprop is used as thermodynamic backend
- Only pure substances are considered
- Gas phase only
- No temperature stratification in the gas phase
- No temperature gradient through vessel wall
- Heat transfer is modelled as simple as possible

These choices make the problem a lot more simple to solve. First of all, the the pure substance Helmholtz energy based equation of state (HEOS) in coolprop offers a lot of convenience in terms of the property pairs/state variables that can be set independently. Using only a single gas phase species also means that component balances is redundant and 2 or 3-phase flash calculations are not required. That being said, the principle used for a single component is more or less the same, even for multicomponent mixtures with potentially more than one phase.

1.2 Getting the software

The source code can be obtained from GitHub either via [git](#) or via the latest tar-ball release. No packaged releases have currently been planned for **pip**, **canda** or similar, since the code is still in alpha stage. This might change in the future. The main branch is located here:

<https://github.com/andr1976/HydDown>

1.3 Requirements

- Python (3.8 - at least python3)
- Numpy
- matplotlib
- Coolprop (6.4.1)
- cerberus

- PyYaml
- pandas

The script is running on Windows 10 x64, with stock python installation from python.org and packages installed using pip. Should run on linux (it does on an Ubuntu image on GitHub) or in any conda environment as well, but I haven't checked.

1.4 Units of measure

The SI unit system is adapted for this project. The following common units are used in the present project and this also applies to the units used in the input files:

Table 1.1: Unit system

Property	Unit	Comment
Temperature	K	° C is used in plots
Pressure	Pa	bar is used in plots
Mass	kg	
Volume	m ³	
Time	s	
Energy	J	
Duty/power	W	
Length	m	
Area	m ²	
Heat flux	W/m ²	
Heat transfer coefficient	W/(m ² K)	
Density	kg/m ³	
Heat capacity	J/(kg K)	

As will be noted when presenting the equations implemented in the code, some of the equations utilise different units than the ones listed in tbl. 1.1. However, it is important to note that unit conversions are built in to the methods implemented, so the user shall not worry about unit conversion.

1.5 Credit

In the making of this document a great deal of material has been sourced (and modified) from a good colleagues M.Sc. thesis (Eriksen and Bjerre 2015), from co-published papers (Bjerre et al. 2017)(Andreasen et al. 2018) and from on-line material published under permissive licenses (with proper citation). Further, the making of this project would not have possible without the awesome CoolProp library (Bell et al. 2014). I am also thankful for enlightning discussions with colleague Jacob Gram Iskov Eriksen (Ramboll Energy, Denmark) and former Ramboll Energy colleague Carsten Stegelmann (ORS Consulting) in relation to vessel depressurisation, nozzle flow and heat transfer considerations.

The present document is typeset using Markdown + pandoc with the Eisvogel template.

1.6 License

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2 Usage

2.1 Basic usage

Run the code as simple as:

```
1 python main.py input.yml
```

where main.py is the main script and input.yml is the input file in Yaml syntax.

The Yaml input file is edited to reflect the system of interest.

2.2 Demos

A few demonstrations of the codes capability are available as `streamlit` apps:

- Vessel depressurisation using simple methods for different gases available at <https://share.streamlit.io/andr1976/hyddown/main>
- Hydrogen pressurisation available at https://share.streamlit.io/andr1976/hyddown/main/streamlit_h2app.py

2.3 Calculation methods

The following methods are implemented:

- Isothermal i.e. constant temperature of the fluid during depressurisation (for a very slow process with a large heat reservoir)
- Isenthalpic/Adiabatic (no heat transfer with surroundings, no work performed by the expanding fluid for depressurisation, no work added from inlet stream)
- Isentropic (no heat transfer with surroundings, PV work performed by the expanding fluid)
- Isenergetic i.e. constant internal energy
- Energy balance. This is the most general case and is based on the first law of thermodynamics applied to a flow process.

For `isothermal/isenthalpic/isentropic/isenergetic` calculations the minimal input required are:

- Initial conditions (pressure, temperature)
- vessel dimensions (ID/length)
- valve parameters (Cd, diameter, backpressure)
- Calculation setup (time step, end time)
- Type of gas

If heat transfer is to be considered the calculation type `energybalance` is required. A few options are possible:

- Fixed U (U-value required, and ambient temperature)
- Fixed Q (Q to be applied to the fluid is required)
- Specified h, the external heat transfer coefficient is provided and either the internal is provided or calculated from assumption of natural convection from a vertical cylinder at high Gr number. Ambient temperature is required.
- Fire (Stefan-Boltzmann equation heat duty)

More elaborate description of the required input for the different calculation types are provided in Sec. 2.7.

2.4 Script

HydDown comes with a script which can be used as a command-line tool to start calculations. If an input filename (path) is given as the first argument, this input file will be used. If no arguments are passed, the script will look for an input file with the name `input.yaml`. The content of the main script is shown below.

```
1 import yaml
2 import sys
3 from hyddown import HydDown
4
5 if __name__ == "__main__":
6     if len(sys.argv) > 1:
7         input_filename = sys.argv[1]
8     else:
9         input_filename = "input.yaml"
10
11     with open(input_filename) as infile:
12         input = yaml.load(infile, Loader=yaml.FullLoader)
13
14
```

```
15     hdown=HydDown(input)
16     hdown.run()
17     hdown.verbose=1
18     hdown.plot()
```

2.5 Module import

To use HydDown simple import the main calculation class `HydDown`.

```
1 from hyddown import HydDown
```

2.6 Input file examples

When using HydDown a dictionary holding all relevant input in order for HydDown to do vessel calculations shall be provided when the class is initialized. One way is to read an input file. For HydDown a Yaml format is chosen, but if JSON is a preference this should also work with the Yaml parser being substituted with a JSON parser.

An example of a minimal file for an isentropic vessel depressurisation (no heat transfer) is shown below

```
1 vessel:
2   length: 1.524
3   diameter: 0.273
4   initial:
5     temperature: 388.0
6     pressure: 15000000.
7     fluid: "N2"
8   calculation:
9     type: "isentropic"
10    time_step: 0.05
11    end_time: 100.
12  valve:
13    flow: "discharge"
14    type: "orifice"
15    diameter: 0.00635
16    discharge_coef: 0.8
17    back_pressure: 101300.
```

A more elaborate example which includes heat transfer and with validation data (some data points dropped for simplicity) included:

```
1 vessel:
2   length: 1.524
```

```

3   diameter: 0.273
4   thickness: 0.025
5   heat_capacity: 500
6   density: 7800.
7   orientation: "vertical"
8   initial:
9     temperature: 288.0
10    pressure: 15000000.
11    fluid: "N2"
12  calculation:
13    type: "energybalance"
14    time_step: 0.05
15    end_time: 100.
16  valve:
17    flow: "discharge"
18    type: "orifice"
19    diameter: 0.00635
20    discharge_coef: 0.8
21    back_pressure: 101300.
22  heat_transfer:
23    type: "specified_h"
24    temp_ambient: 288.
25    h_outer: 5
26    h_inner: 'calc'
27  validation:
28    temperature:
29      gas_high:
30        time: [0.050285, ... , 99.994]
31        temp: [288.93, ... ,241.29]
32      gas_low:
33        time: [0.32393, ... , 100.11]
34        temp: [288.67, ... ,215.28]
35      wall_low:
36        time: [0.32276, ... , 100.08]
37        temp: [288.93, ... ,281.72]
38      wall_high:
39        time: [0.049115, ... ,100.06]
40        temp: [289.18, ... ,286.09]
41    pressure:
42      time: [0.28869, ... , 98.367]
43      pres: [150.02, ... ,1.7204]

```

2.7 Input fields and hierachy

In the following the full hierachy of input for the different calculation types is summarised.

At the top level the following fields are accepted, with the last being optional and the second last dependant on calculation type:

```
1 initial: mandatory
2 vessel: mandatory
3 calculation: mandatory
4 valve: mandatory
5 heat_transfer: depends on calculation type
6 validation: optional
```

2.7.1 Calculation

The subfields under `calculation`, with value formats and options are:

```
1 calculation:
2   type: "isothermal", "isentropic", "isenthalpic", "constantU", "
      energybalance"
3   time_step: number
4   end_time: number
```

The simulation end time is specified as well as the fixed time step used in the integration of the differential equations to be solved. The four main calculation types are shown as well.

2.7.2 Vessel

```
1 vessel:
2   length: number, mandatory
3   diameter: number, mandatory
4   thickness: number, required when heat transfer is calculated
5   heat_capacity: number, required when heat transfer is calculated
6   density: number, required when heat transfer is calculated
7   orientation: string, required when heat transfer is calculated
```

2.7.3 Initial

```
1 initial:
2   temperature: number, mandatory
3   pressure: number, mandatory
4   fluid: string, mandatory , e.g. "N2", "H2"
```

2.7.4 Valve

The `valve` field determines the mode of the process is it for depressurisation/discharge from the vessel using the value `discharge` or if mass flow is entering the vessel, filling it up/increasign pressure with

the value `filling`.

Different types of mass flow devices can be specified:

- Restriction `orifice`
- Relief valve/pressure safety valve (only for `discharge` not `filling`)
- Control valve (`controlvalve`)
- A specified mass flow (`mdot`)

For the physical devices a `back_pressure` is required for the flow calculations. The value of the `back_pressure` **is also used to specify the reservoir pressure when the vessel is filled**. See also Sec. 3.2 for details about the calculation of flow rates.

```
1 valve:
2   flow: string, mandatory "discharge" or "filling"
3   type: string, mandatory "orifice", "controlvalve", "psv", "mdot"
4   back_pressure: number, required for type "orifice", "!controlvalve"
      and "psv"
5   diameter: number, required for "orifice" and "psv"
6   discharge_coef: number, required for "orifice" and "psv"
7   Cv: number, required for "control_valve"!
```

2.7.5 Heat transfer

For more information about the actual estimation of heat transfer see also Sec. 3.3. For the case of `fire` heat input predetermined parameters for the Stefan-Boltzmann fire equation are used to calculate different background heat loads cf. Tbl. 2.1

Table 2.1: Fire heat loads

Source	Fire type	Heat load (kW/m^2)
API521	Pool	60
API521	Jet	100
Scandpower	Pool	100
Scandpower	Jet	100

```
1 heat_transfer:
2   type: string, mandatory, "specified_h", "specified_Q", "specified_U",
      "s-b"
3   temp_ambient: number, required for type "specified_h", "specified_U"
```

```
4  h_outer: number, required for type "specified_h"
5  h_inner: number or 'calc', required for type "specified_h"
6  U_fix: number, required for type "specified_U"
7  Q_fix: number, required for type "specified_Q"
8  fire: string, required for type "s-b", 'api_pool', 'api_jet', '
      scandpower_pool', 'scandpower_jet'
9  D_throat: number, required for flow type "filling", set to vessel ID
      as a starting point
```

2.7.6 Validation

In order to plot measured data against simulated data the field `validation` is included.

The following arrays (one or more) are supported in the sub-field `temperature`:

- `gas_high`: highest measured values of the bulk gas
- `gas_low`: lowest measured values of the bulk gas
- `gas_mean`: average measured values of the bulk gas

- `wall_mean`: average measured values of the vessel (inner) wall
- `wall_high` i.e. highest measured values of the vessel (inner) wall
- `wall_low` i.e. lowest measured values of the vessel (inner) wall

For each of the above fields arrays for `time` and `temp` shall be supplied with matching length.

A field for measured vessel pressure is also possible, where `time` and `pres` shall be identical length arrays. See also example below.

```
1  validation:
2    temperature:
3      gas_high:
4        time: [0.050285, ... , 99.994]
5        temp: [288.93, ... , 241.29]
6      gas_low:
7        time: [0.32393, ... , 100.11]
8        temp: [288.67, ... , 215.28]
9      wall_low:
10       time: [0.32276, ... , 100.08]
11       temp: [288.93, ... , 281.72]
12     wall_high:
13       time: [0.049115, ... , 100.06]
14       temp: [289.18, ... , 286.09]
15     pressure:
16       time: [0.28869, ... , 98.367]
17       pres: [150.02, ... , 1.7204]
```


3 Theory

In this chapter the basic theory and governing equations for the model implementation in HydDown is presented. The following main topics are covered:

- thermodynamics
- mass transfer
- heat transfer

3.1 Thermodynamics

3.1.1 Equation of state

The equation of state used by HydDown is the Helmholtz energy formulation as implemented in CoolProp (Bell et al. 2014). Most of the text in the present section has been sourced from the CoolProp documentation to be as accurate and true to the source as possible. The Helmholtz energy formulation is a convenient construction of the equation of state because all the thermodynamic properties of interest can be obtained directly from partial derivatives of the Helmholtz energy.

It should be noted that the EOS are typically valid over the entire range of the fluid, from subcooled liquid to superheated vapor, to supercritical fluid. In general, the EOS are based on non-dimensional terms δ and τ , where these terms are defined by

$$\delta = \rho/\rho_c$$

$$\tau = T_c/T$$

where ρ_c and T_c are the critical density of the fluid if it is a pure fluid. For pseudo-pure mixtures, the critical point is typically not used as the reducing state point, and often the maximum condensing temperature on the saturation curve is used instead.

The non-dimensional Helmholtz energy of the fluid is given by

$$\alpha = \alpha^0 + \alpha^r$$

where α^0 is the ideal-gas contribution to the Helmholtz energy, and α^r is the residual Helmholtz energy contribution which accounts for non-ideal behavior. For a given set of δ and τ , each of the terms α^0 and α^r are known. The exact form of the Helmholtz energy terms is fluid dependent, but a relatively simple example is that of Nitrogen, which has the ideal-gas Helmholtz energy of

$$\alpha^0 = \ln \delta + a_1 \ln \tau + a_2 + a_3 \tau + a_4 \tau^{-1} + a_5 \tau^{-2} + a_6 \tau^{-3} + a_7 \ln[1 - \exp(-a_8 \tau)]$$

and the non-dimensional residual Helmholtz energy of

$$\alpha^r = \sum_{k=1}^6 N_k \delta^{i_k} \tau^{j_k} + \sum_{k=7}^{32} N_k \delta^{i_k} \tau^{j_k} \exp(-\delta^{l_k}) + \sum_{k=33}^{36} N_k \delta^{i_k} \tau^{j_k} \exp(-\phi_k(\delta - 1)^2 - \beta_k(\tau - \gamma_k)^2)$$

and all the terms other than δ and τ are fluid-dependent correlation parameters.

The other thermodynamic parameters can then be obtained through analytic derivatives of the Helmholtz energy terms. For instance, the pressure is given by

$$p = \rho RT \left[1 + \delta \left(\frac{\partial \alpha^r}{\partial \delta} \right)_{\tau} \right]$$

and the specific internal energy by

$$\frac{u}{RT} = \tau \left[\left(\frac{\partial \alpha^0}{\partial \tau} \right)_{\delta} + \left(\frac{\partial \alpha^r}{\partial \tau} \right)_{\delta} \right]$$

and the specific enthalpy by

$$\frac{h}{RT} = \tau \left[\left(\frac{\partial \alpha^0}{\partial \tau} \right)_{\delta} + \left(\frac{\partial \alpha^r}{\partial \tau} \right)_{\delta} \right] + \delta \left(\frac{\partial \alpha^r}{\partial \delta} \right)_{\tau} + 1$$

which can also be written as

$$\frac{h}{RT} = \frac{u}{RT} + \frac{p}{\rho RT}$$

The specific entropy is given by

$$\frac{s}{R} = \tau \left[\left(\frac{\partial \alpha^0}{\partial \tau} \right)_{\delta} + \left(\frac{\partial \alpha^r}{\partial \tau} \right)_{\delta} \right] - \alpha^0 - \alpha^r$$

and the specific heats at constant volume and constant pressure respectively are given by

$$\frac{c_v}{R} = -\tau^2 \left[\left(\frac{\partial^2 \alpha^0}{\partial \tau^2} \right)_{\delta} + \left(\frac{\partial^2 \alpha^r}{\partial \tau^2} \right)_{\delta} \right]$$

$$\frac{c_p}{R} = \frac{c_v}{R} + \frac{\left[1 + \delta \left(\frac{\partial \alpha^r}{\partial \delta} \right)_{\tau} - \delta \tau \left(\frac{\partial^2 \alpha^r}{\partial \delta \partial \tau} \right) \right]^2}{\left[1 + 2\delta \left(\frac{\partial \alpha^r}{\partial \delta} \right)_{\tau} + \delta^2 \left(\frac{\partial^2 \alpha^r}{\partial \delta^2} \right)_{\tau} \right]}$$

The EOS is set up with temperature and density as the two independent properties, but often other inputs are known, most often temperature and pressure because they can be directly measured. As a result, if the density is desired for a known temperature and pressure, it can be obtained iteratively.

3.1.2 First law for flow process

The control volume sketched in Fig. 3.1, separated from the surrounding by a control surface, is used as a basis for the analysis of an open thermodynamic system with flowing streams (fs) in and out, according to (Smith, Van Ness, and Abbott 1996)

A general mass balance or continuity equation can be written:

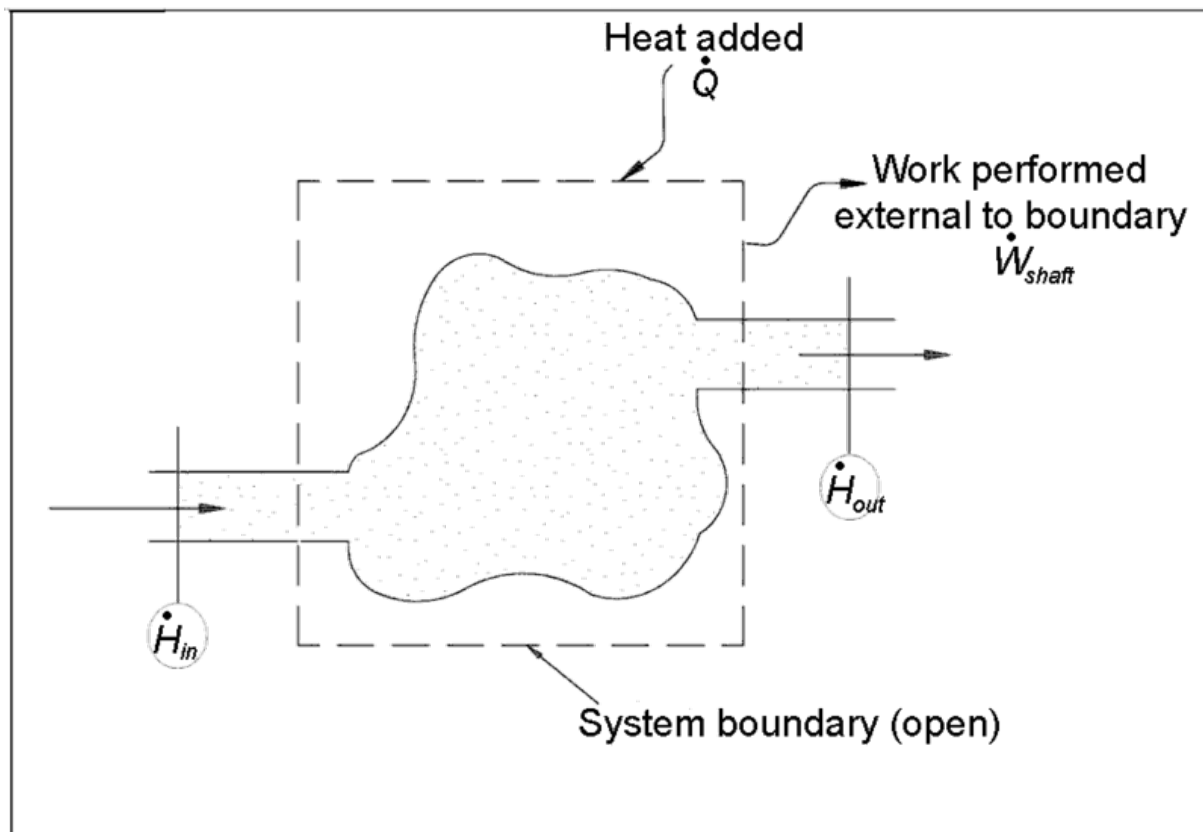


Figure 3.1: Control volume with one entrance and one exit. The image has been sourced from (WikiMedia 2008,).

$$\frac{m_{cv}}{dt} + \Delta(\dot{m})_{fs} = 0 \quad (3.1)$$

The first term is the accumulation term i.e. the rate of change of the mass inside the control volume, m_{cv} , and the Δ in the second term represents the difference between the outflow and the inflow

$$\Delta(\dot{m})_{fs} = \dot{m}_2 - \dot{m}_1$$

An energy balance for the control volume, with the first law of thermodynamics applied, needs to account for all the energy modes with can cross the control surface. Each stream has a total energy

$$U + \frac{1}{2}u^2 + zg$$

where the first term is the specific internal energy, the second term is the kinetic energy and the last

term is the potential energy. The rate at which each stream transports energy in or out of the control volume is given by

$$\dot{m}(U + \frac{1}{2}u^2 + zg)$$

and in total

$$\Delta \left[\dot{m}(U + \frac{1}{2}u^2 + zg) \right]_{fs}$$

Further work (not to be confused with shaft work) is also associated with each stream in order to move the stream into or out from the control volume (one can think of a hypothetical piston pushing the fluid at constant pressure), and the work is equal to PV on the basis of the specific fluid volume. The work rate for each stream is

$$\dot{m}(PV)$$

and in total

$$\Delta [\dot{m}(PV)]_{fs}$$

Further, heat may be transferred to (or from) the control volume at a rate \dot{Q} and shaft work may be applied, \dot{W}_{shaft} . Combining all this with the accumulation term given by the change in total internal energy the following general energy balance can be written

$$\frac{d(mU)_{cv}}{dt} + \Delta \left[\dot{m}(U + \frac{1}{2}u^2 + zg) \right]_{fs} + \Delta [\dot{m}(PV)]_{fs} = \dot{Q} + \dot{W}_{shaft}$$

Applying the relation $H = U + PV$, setting $\dot{W}_{shaft} = 0$ since no shaft work is applied to or extracted from the vessel, and assuming that kinetic and potential energy changes can be omitted the energy balance simplifies to:

$$\frac{d(mU)_{cv}}{dt} + \Delta [\dot{m}H]_{fs} = \dot{Q}$$

The equation can be further simplified if only a single port acting as either inlet or outlet is present

$$\frac{d(mU)_{cv}}{dt} + \dot{m}H = \dot{Q} \tag{3.2}$$

where the sign of \dot{m} determines if the control volume is either emptied or filled. The continuity equation Eq. 3.1 and the energy balance Eq. 3.2 combined with the equation of state are the key equations that shall be solved/integrated in order to calculate the change in temperature and pressure as a function of time.

3.2 Flow devices

3.2.1 Restriction Orifice

When a fluid flows through a restriction or opening such as an orifice, the velocity will be affected by conditions upstream and downstream. If the upstream pressure is high enough, relative to the downstream pressure, the velocity will reach the speed of sound ($Ma = 1$) and the flow rate obtained will be the critical flow rate. The maximum downstream pressure for the flow to still be sonic ($Ma = 1$), is when $P_d = P_c$. The ratio of the critical and upstream pressure is defined by equation Eq. 3.3.

$$\frac{P_c}{P_u} = \left(\frac{2}{k+1} \right)^{\frac{k}{k-1}} \quad (3.3)$$

- P_c is the critical pressure. [kPa]
- P_d is the downstream pressure. [kPa]
- P_u is the upstream pressure. [kPa]
- k is the isentropic coefficient, approximated by the ideal gas heat capacity ratio C_p/C_v . [-]

In order to calculate the mass flow rate through an orifice equations are used based on literature from the Committee for the Prevention of Disasters (Bosch and Weterings 2005).

To account for the difference in choked and non-choked flow a set limit pressure is introduced as in equation Eq. 3.4. If the downstream pressure, P_{down} , is below the pressure limit, P_{limit} , then the flow is choked, and the pressure used, P_{used} , in equation Eq. 3.5 should be the pressure limit, P_{limit} . Otherwise if the downstream pressure, P_{down} , is greater than or equal to the pressure limit, P_{limit} , the flow is no longer choked and the pressure used should be the downstream pressure, P_{down} (Bosch and Weterings 2005).

$$P_{limit} = P_{up} \cdot \left(\frac{2}{k+1} \right)^{\frac{k}{k-1}} \quad (3.4)$$

$$\dot{m}_{flow} = C_d \cdot A \cdot \sqrt{\left(\frac{2k}{k-1} \right) \cdot P_{up} \cdot \rho \cdot \left(\frac{P_{used}}{P_{up}} \right)^{\frac{2}{k}} \left(1 - \left(\frac{P_{used}}{P_{up}} \right)^{\frac{k-1}{k}} \right)} \quad (3.5)$$

- ρ is the density of the gas upstream. $[kg/m^3]$
- P_{limit} is the pressure limit of the upstream absolute pressure. $[bara]$
- P_{up} is the absolute pressure upstream of the orifice. $[bara]$
- k is the ratio of the heat capacities at constant pressure, C_p , and at constant volume, C_v .
- P_{down} is the absolute pressure downstream of the orifice. $[bara]$
- P_{used} is the pressure used in the mass flow equation based on choked or non-choked conditions. $[bara]$
- \dot{m}_{flow} is the mass flow through the orifice. $[kg/s]$
- C_d is the discharge coefficient of the orifice opening. $[-]$
- A is the cross sectional area of the orifice. $[m^2]$

3.2.2 Pressure safety valve / Relief valve

A PSV / relief valve is a mechanical device actuated by the static pressure in the vessel and a conventional PSV is often used for gas/vapor systems. A conventional PSV is a spring-loaded device which will activate at a predetermined opening pressure, and relieve the vessel pressure until a given reseal pressure has been reached. Both the opening pressure and the reset pressure is above the vessel operating pressure, and the PSV will remain closed until the pressure inside the vessel increases to the opening pressure. The operation of a conventional spring-loaded PSV is based on a force balance. A conventional PSV can be seen in Fig. 3.2. A spring exerts a force on a disc blocking the inlet of the PSV. When the pressure inside the vessels reaches the opening pressure, the force exerted on the disc by the gas, will be larger than the force exerted by the spring and the PSV will open and allow the gas to flow out of the vessels. The flow of gas out of the vessels will lower the pressure and thereby also the force exerted on the disc. When the pressure in the vessels is reduced to the reset pressure, the PSV will close and the disc will again hinder the gas flow.

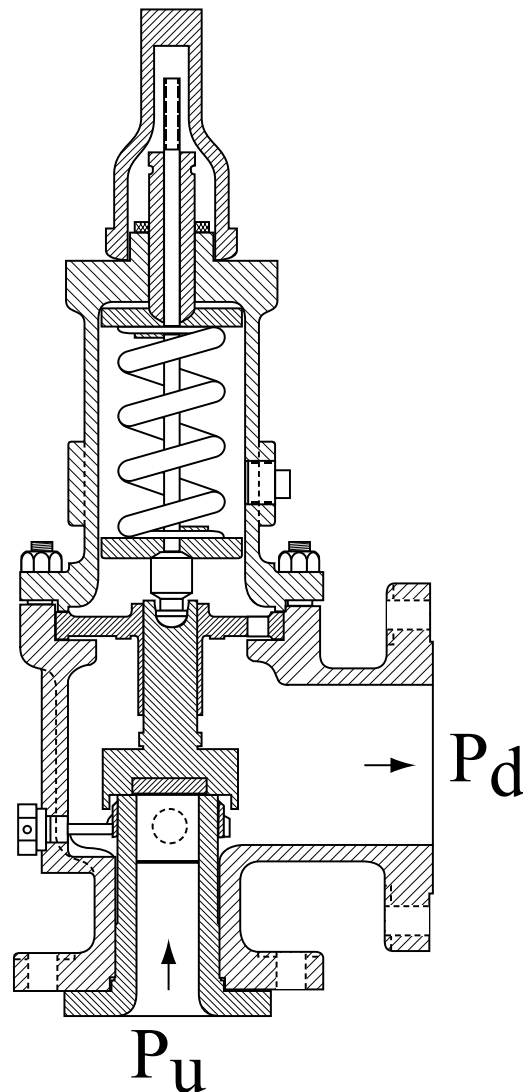


Figure 3.2: Conventional/pop action PSV adapted from (Eriksen and Bjerre 2015) and (API 2014a)

The relief valve model implemented in HydDown is the API 520 equations (API 2014a) for gas relief for both sonic/critical as well as subcritical flow. No corrections factors are implemented in HydDown.

For sonic flow (critical flow), as indicated in equation, the mass flow thorough the PSV can be determined by equation eq. 3.6.

$$W = \frac{AC \cdot K_d \cdot K_b \cdot K_c \cdot P_1}{\sqrt{\frac{T \cdot Z}{M}}} \quad (3.6)$$

- A is the effective discharge area. [mm²]
- W is the mass flow through the device. [kg/h]

- C is a coefficient, as a function of k , as defined in equation eq. 3.7.
- K_d , K_b , and K_c are correction factors.
- P_1 is the allowable upstream absolute pressure. [kPa]
- T is the temperature of the inlet gas at relieving conditions. [K]
- M is the molecular mass of the gas at relieving conditions. [kg/kmol]
- Z is the compressibility factor for the gas.

K_d is the effective coefficient of discharge, with a typical value of 0.975, for an installed PSV. K_b is a back-pressure correction factor between 0 and 1, assumed to be 1. K_c is a correction factor used when a rupture disk is installed upstream, otherwise it is 1. In the present implementation a value of 1 is assumed.

$$C = 0.03948 \sqrt{k \left(\frac{2}{k+1} \right)^{\left(\frac{k+1}{k-1} \right)}} \quad (3.7)$$

For subsonic flow (subcritical flow), the effective discharge area of the PSV is determined by equation eq. 3.8.

$$W = \frac{A \cdot F_2 \cdot K_d \cdot K_c}{17.9 \sqrt{\frac{T \cdot Z}{M \cdot P_1 \cdot (P_1 - P_2)}}} \quad (3.8)$$

F_2 is the coefficient of subcritical flow which can be determined from eq. 3.9.

$$F_2 = \sqrt{\left(\frac{k}{k-1} \right) r^{\left(\frac{2}{k} \right)} \left(\frac{1 - r^{\left(\frac{k-1}{k} \right)}}{1 - r} \right)} \quad (3.9)$$

where r is the ratio of backpressure to upstream relieving pressure, P_2/P_1 .

When modelling a pop action PSV/relief valve under dynamic conditions, the valve will go from closed to fully open in a short period of time when the set pressure, P_{set} , is reached. The pop action is illustrated in Fig. 3.3 which shows the opening and closing hysteresis of the PSV as a function of pressure. In order to close the pressure shall be reduced below the reseal pressure.

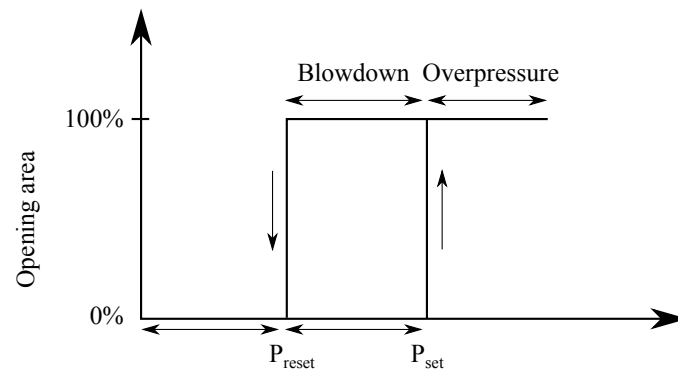


Figure 3.3: Relief valve hysteresis adapted from (Eriksen and Bjerre 2015)

When specifying PSV's it is comon to use standrd API sizes as shown in tbl. 3.1

Table 3.1: Standard PSV orifice sizes according to API

Size	Area [in ²]	Area [m ²]
D	0.110	$7.09676 \cdot 10^{-5}$
E	0.196	$1.26451 \cdot 10^{-4}$
F	0.307	$1.98064 \cdot 10^{-4}$
G	0.503	$3.24515 \cdot 10^{-4}$
H	0.785	$5.06450 \cdot 10^{-4}$
J	1.287	$8.30320 \cdot 10^{-4}$
K	1.838	$1.18580 \cdot 10^{-3}$
L	2.853	$1.84064 \cdot 10^{-3}$
M	3.600	$2.32257 \cdot 10^{-3}$
N	4.340	$2.79999 \cdot 10^{-3}$
P	6.380	$4.11612 \cdot 10^{-3}$
Q	11.050	$7.12901 \cdot 10^{-3}$
R	16.000	$1.03225 \cdot 10^{-2}$
T	26.000	$1.67741 \cdot 10^{-2}$

3.2.3 Control Valve

For calculating the mass flow through a control valve the ANSI/ISA (Borden 1998)(ISA 1995) methodology also described in IEC 60534 (IEC 2011) is applied.

The flow model for a compressible fluid in the turbulent regime is

$$W = CN_6 F_P Y \sqrt{x_{sizing} p_1 \rho_1}$$

or equivalent

$$W = CN_8 F_P Y \sqrt{\frac{x_{sizing} M}{T_1 Z_1}}$$

- C is the flow coefficient (C_v or K_v)
- N_8 is a unit specific constant, 94.8 for C_v and bar as pressure unit
- F_P is a piping geometry factor [-]
- Y is the expansion factor [-]
- x_{sizing} is the pressure drop used for sizing[-]
- p_1 is the upstream pressure [bar]
- ρ_1 is the upstream density [kg/m³]
- M is the molecular weight [kg/kmol]
- T_1 is the upstream temperature [K]
- Z_1 is the upstream compressibility [-]

In HydDown the piping geometry factor is not yet implemented and asumed to be 1. The pressure drop ratio x_{sizing} used for sizing is determined as the lesser of the actual pressure drop ratio, x , and the choked pressure drop ratio x_{choked} . The actual pressure drop ratio is given by:

$$x \frac{\Delta p}{p_1}$$

The pressure drop ratio at which flow no longer increases with increased value in pressure drop ratio, is the choked pressure drop ratio, given by the following equation

$$x_{choked} = F_\gamma x_{TP}$$

The factor x_T is based on air near atmospheric pressure as the flowing fluid with a specific heat ratio of 1.40. If the specific heat ratio for the flowing fluid is not 1.40, the factor F_γ is used to adjust x_T . Use the following equation to calculate the specific heat ratio factor:

$$F_\gamma = \frac{\gamma}{1.4}$$

where γ is the ideal gas C_p/C_v . It should be noted that the above equation has been derived from perfect gas behaviour and extension of an orifice model with γ in the range of 1.08 to 1.65. If used outside the assumptions flow calculations may become inaccurate.

The expansion factor Y accounts for the change in density as the fluid passes from the valve inlet to the vena contracta. It also accounts for the change in the vena contracta area as the pressure differential is varied.

$$Y = 1 - \frac{x_{sizing}}{3x_{choked}}$$

3.3 Heat transfer

3.3.1 Natural convection

Experiments have indicated that the internal heat transfer mechanism for a vessel subject to depressurisation can be well approximated by that of natural convection as found from measured Nusselt numbers being well correlated with Rayleigh number, with no apparent improvement in model performance by inclusion of the Reynold number in the model (Woodfield, Monde, and Mitsutake 2007).

To determine the heat transfer for the gas-wall interface, the following is applied cf. equation Eq. 3.10.

$$\frac{dQ}{dt} = \dot{Q} = hA(T_s - T_{gas}) \quad (3.10)$$

- dQ is the change in thermal energy due to convective heat transfer. [J]
- dt is the change in time during the heat transfer. [s]
- h is the convective heat transfer. [$\text{W}/\text{m}^2 \cdot \text{K}$]
- A is the area normal to the direction of the heat transfer. [m^2]
- T_s is the inner surface temperature of the geometry. [K]
- T_{gas} is the temperature of the bulk gas inside the vessel. [K]

The convective heat transfer will need to be estimated for the the gas-wall interface, by the use of empirical relations for the Nusselt number. The Nusselt number describes the ratio of convective heat transfer to conductive heat transfer, normal to a surface area, as given in equation eq. 3.11.

$$Nu = \frac{hL}{k} \quad (3.11)$$

- Nu is the Nusselt number. [-]
- h is the convective heat transfer. [$\text{W}/\text{m}^2\cdot\text{K}$]
- L is a characteristic length of the geometry. [m]
- k is the thermal conductivity of the gas. [$\text{W}/\text{m}\cdot\text{K}$]

The characteristic length L used is the height of the gas volume i.e. the length of a vertical vessel or the diameter of a horizontal vessel.

The empirical correlations used to calculate the Nusselt number of the gas-wall interface is a function of the Rayleigh number, which can be defined by the Grashof number and Prandtl number, as in equation Eq. 3.12.

$$Ra = Gr \cdot Pr \quad (3.12)$$

- Ra is the Rayleigh number. [-]
- Gr is the Grashof number. [-]
- Pr is the Prandtl number. [-]

The Grashof number is a dimensionless number which approximates the ratio of the buoyancy forces to viscous forces, as given in equation [Eq. 3.13]. The Prandtl number is a dimensionless number defined as the ratio of the momentum diffusivity to thermal diffusivity, as given in equation Eq. 3.14.

$$Gr = \frac{\beta g \rho^2 L^3 \Delta T}{\mu^2} \quad (3.13)$$

$$Pr = \frac{c_p \mu}{k} \quad (3.14)$$

- β is the coefficient of volume expansion. [$1/\text{K}$]
- g is the standard acceleration of gravity. [m/s^2]
- ρ is the gas density. [kg/m^3]
- L is the characteristic length. [m]
- ΔT is the temperature difference of the surface and gas. [K]
- μ is the dynamic viscosity. [$\text{kg}/\text{m}\cdot\text{s}$]
- c_p is the heat capacity of gas. [$\text{J}/\text{kg}\cdot\text{K}$]
- k is the thermal conductivity of gas. [$\text{J}/\text{m}\cdot\text{K}$]

It is important to note that the properties in the above equations shall be evaluated at the fluid film temperature which can be approximated by the average of the the fluid bulk temperature and the vessel wall temperature (Geankoplis 1993).

The mechanism of heat transfer on the outside of the vessel at ambient conditions is similar to the above. In HydDown the outside heat transfer is not modelled currently. A heat transfer coefficient shall be provided.

3.3.2 Mixed convection

Experiments have indicated that the internal heat transfer mechanism for a vessel subject to filling can be well approximated by that of combined natural convection and forced convection as found from measured Nusselt numbers being well correlated with Rayleigh and Reynolds number (Woodfield, Monde, and Mitsutake 2007).

For mixed convection the effective Nusselt number, Nu , can be approximated by

$$Nu = (Nu_{forced}^n + Nu_{natural}^n)^{\frac{1}{n}}$$

During charging with different gases (H₂, N₂ and Argon), Woodfield *et al.* demonstrated that in order to provide a good fit to the experimentally determined Nusselt number a correlation based on both Reynolds and Rayleigh number was necessary. They found a good fit with the following formula (n=1)

$$Nu = Nu_{forced} + Nu_{natural} = 0.56Re_d^{0.67} + 0.104Ra_H^{0.352}$$

3.3.3 Conduction

For accurate prediction of the outer and especially the inner wall temperature for correct estimation of internal convective heat transfer and the average materials temperature, the general equation of 1-D unsteady heat transfer shall be solved:

$$\frac{\delta T}{\delta t} = \frac{k}{C_p} \frac{\delta^2 T}{\delta x^2}$$

- T is temperature
- x is the spatial (1-D) coordinate
- k is the thermal conductivity
- C_p is the heat capacity

Here it is written in Cartesian coordinates, but for most applications to pressure equipment, cylindrical coordinates are applicable, at least for the shell. To be solved the initial values and boundary values must be specified. In its present state HydDown does not include the unsteady heat transfer model i.e. the assumption is that the temperature from outer to inner surface is uniform and equal to the

average temperature. This is obviously a crude approximation, but might be justified depending in the Biot number:

$$Bi = \frac{hL}{k}$$

The Biot number is a simple measure of the ratio of the heat transfer resistances at the surface of a body to the inside of a body. The ratio gives an indication to which extent the temperature will vary in space (gradient) when the body is subject to a displacement in temperature at the surface boundary layer. Striednig *et al.* (Striednig et al. 2014) concluded that for a type I (steel) cylinder the Biot number was approx. 0.03 and hence the error in assuming a uniform temperature in the vessel wall was low.

With a typical thermal conductivity of 45 W/mK for steel and a heat transfer coefficient up to $600 \text{ W/m}^2\text{K}$ (Woodfield, Monde, and Mitsutake 2007) the Biot number for a vessel with a wall thickness of 2 cm is 0.27. This is significant higher than approximated by (Striednig et al. 2014). Anyway, the Biot number is lower than 1, and the assumption of a uniform temperature is reasonable. However, for increased wall thickness, and/or for different materials with lower thermal conductivity the error may grow to an unacceptable level.

3.3.4 Fire heat loads

The heat transfer from the flame to the shell is modelled using the recommended approach from Scandpower (Hekkelstrand and Skulstad 2004) and API (API 2014b). The heat transfer from the flame to the vessel shell is divided into radiation, convection and reradiation as seen in equation Eq. 3.15.

$$q_f = \underbrace{\alpha_s \cdot \varepsilon_f \cdot \sigma \cdot T_f^4}_{\text{Radiation}} + \underbrace{h_f \cdot (T_f - T_s(t))}_{\text{Convection}} - \underbrace{\varepsilon_s \cdot \sigma \cdot T_s(t)^4}_{\text{Reradiation}} \quad (3.15)$$

- q_f is the flame heat flux. [W/m^2]
- α_s is the vessel surface absorptivity. [-]
- ε_f is the flame emissivity. [-]
- σ is the Stefan-Boltzmann constant, $\sigma = 5.67 \cdot 10^{-8} \text{ [W/m}^2 \cdot \text{K}^4]$
- T_f is the flame temperature. [K]
- h_f is the convection heat transfer coefficient between the flame and the surface. [$\text{W/m}^2 \cdot \text{K}$]
- $T_s(t)$ is the time dependent surface temperature. [K]
- ε_s is the surface emissivity. [-]

This model assumes that the pressure vessel is fully engulfed by the flame. This means that the view factor for the radiation is unity and is therefore not taken into consideration. The convective heat

transfer coefficients for a jet fire and a pool fire, and recommended values for the emissivity and absorptivity, are given by Scandpower as (Hekkelstrand and Skulstad 2004)

- $h_{jet\ fire} = 100\ [W/m^2 \cdot K]$
- $h_{pool\ fire} = 30\ [W/m^2 \cdot K]$
- $\alpha_s = 0.85$
- $\varepsilon_s = 0.85$
- $\varepsilon_f = 1.0$ (optical thick flames, thickness $> 1\ m$)

The flame temperature is found by solving equation Eq. 3.16 for the incident heat flux in relation to the ambient conditions. The flame temperature is kept constant throughout the simulation.

$$q_{total} = \sigma \cdot T_f^4 + h_f \cdot (T_f - T_{amb}) \quad (3.16)$$

- q_{total} is the incident flame heat flux as given in table Tbl. 3.2. $[W/m^2]$
- T_{amb} is the ambient temperature $\approx 293\ K$ ($20^\circ\ C$)

The heat flux used to calculate the flame temperature is given in table tbl. 3.2.

Table 3.2: Incident heat fluxes for various fire scenarios given by Scandpower (Hekkelstrand and Skulstad 2004)

	Small jet fire $[kW/m^2]$	Large jet fire $[kW/m^2]$	Pool fire $[kW/m^2]$
Peak heat load	250	350	150
Background heat load	0	100	100

3.4 Model implementation

A simple (naive) explicit Euler scheme is implemented to integrate the mass balance over time, with the mass rate being calculated from an orifice/valve equation as described in Sec. 3.2:

$$m_{cv}(i+1) = m_{cv}(i) + \dot{m}(i)\Delta t \quad (3.17)$$

$$\dot{m}(i) = f(P, T,)$$

For each step, the mass relief/ left in the vessel is known. Since the volume is fixed the mass density is

directly given.

For the calculation methods (isentropic, isenthalpic, isenergetic etc), Coolprop allows specifying density and either H, S or U directly - this is very handy and normally only TP, PH, TS property pairs are implemented, and you would need to code a second loop to make it into an UV, VH or SV calculation.

In the following the high level steps in the solution procedure is outlined for the different calculation types. To illustrate calls to the equation of state in `coolprop` the following notation is used $EOS(D, T)$ corresponding to calculation of the state at known density and temperature as an example.

3.4.1 Isothermal process

For an isothermal process the solution procedure for each calculation step is the following, with the mass in the control volume being calculated from Eq. 3.17:

$$\begin{aligned}T(i+1) &= T(i) \\ D(i+1) &= \frac{m_{cv}(i+1)}{V} \\ P(i+1) &= EOS(D(i+1), T(i+1))\end{aligned}$$

3.4.2 Isenthalpic process

For an isenthalpic process

$$\begin{aligned}H(i+1) &= H(i) \\ D(i+1) &= \frac{m_{cv}(i+1)}{V} \\ P(i+1) &= EOS(D(i+1), H(i+1)) \\ T(i+1) &= EOS(D(i+1), H(i+1))\end{aligned}$$

3.4.3 Isentropic process

For an isentropic process

$$\begin{aligned}S(i+1) &= S(i) \\ D(i+1) &= \frac{m_{cv}(i+1)}{V}\end{aligned}$$

$$P(i+1) = EOS(D(i+1), S(i+1))$$

$$T(i+1) = EOS(D(i+1), S(i+1))$$

3.4.4 Isenergetic process

For an isenergetic process

$$U(i+1) = U(i)$$

$$D(i+1) = \frac{m_{cv}(i+1)}{V}$$

$$P(i+1) = EOS(D(i+1), U(i+1))$$

$$T(i+1) = EOS(D(i+1), U(i+1))$$

3.4.5 Energy balance

The general first law applied to a flow process as outlined in Sec. 3.1.2 subject to an explicit Euler scheme is:

$$D(i+1) = \frac{m_{cv}(i+1)}{V}$$

$$U_{cv}(i+1) = \frac{m_{cv}(i)U_{cv}(i) - (\dot{m}(i)H(i) + \dot{Q}(i)) \Delta t}{m_{cv}(i+1)} \quad (3.18)$$

The above assumes that mass flow is positive when leaving the control volume and heat rate is positive when transferred to the control volume. $H(i)$ is the specific enthalpy of the fluid in the control volume for a discharging process and it is equal to the energy of the entering stream for a filling process. The heat rate is calculated as outlined in [#Sec:heat].

For the vessel wall a simple heat balance is also made:

$$T_{wall}(i+1) = T_{wall}(i) + \frac{\dot{Q}_{outer} - \dot{Q}_{inner}}{c_p m_{vessel}} \Delta t$$

where \dot{Q}_{outer} is the convective heat transfer to/from the ambient surroundings from the outer surface of the vessel, with the a possitive value when heat is transferred from surroundings to the vessel wall. This is either a fixed heat transfer coefficient (for now) with a specified ambient temperature, or a

calculated fire heat load. \dot{Q}_{inner} is the internal heat transfer, either a fixed number or calculated as natural convection (for discharge) or mixed natural and forced convection (filling).

If a fixed \dot{Q} is provided or a fixed overall heat transfer coefficient is provided the vessel wall temperature heat balance is not solved.

The remaining steps are update of temperature and pressure

$$P(i+1) = EOS(D(i+1), U(i+1))$$

$$T(i+1) = EOS(D(i+1), U(i+1))$$

4 Validation

The code is provided as-is. However, comparisons have been made to a few experiments from the literature.

The following gases and modes are considered:

- High pressure hydrogen discharge
- High pressure hydrogen filling
- High pressure nitrogen discharge

4.1 Hydrogen discharge

Calculations with HydDown for high pressure hydrogen vessel discharge have been compared to three experiments from (Byrnes, Reid, and Ruccia 1964). Byrnes *et al.* (Byrnes, Reid, and Ruccia 1964) have reported a number of rapid depressurisation experiments using both nitrogen and hydrogen, with blowdown times being varied between 14 seconds to 33 minutes. The vessel used was a type “K” gas cylinder with a volume of 0.0499 m^3 , an ID of 8.56 inches and shell length of 55 inches. The vessel is placed in a vertical position. A control valve is mounted on the neck of the cylinder with a fixed opening percentage during each experiment, set to provide the required depressurisation rate. It is mentioned that the gas cylinder is placed inside a weather proof cabinet with perlite insulation in order to emulate adiabatic conditions. However the simulation run here are with a finite heat transfer coefficient on the outside of the cylinder. In their paper detailed pressure and temperature traces are provided for three experimental runs: 7, 8, and 9, with depressurisation from 2,000 psia in 30 seconds, 480 seconds and 14 seconds, respectively. These experiments are simulated in Fig. 4.1, Fig. 4.2 and Fig. 4.3.

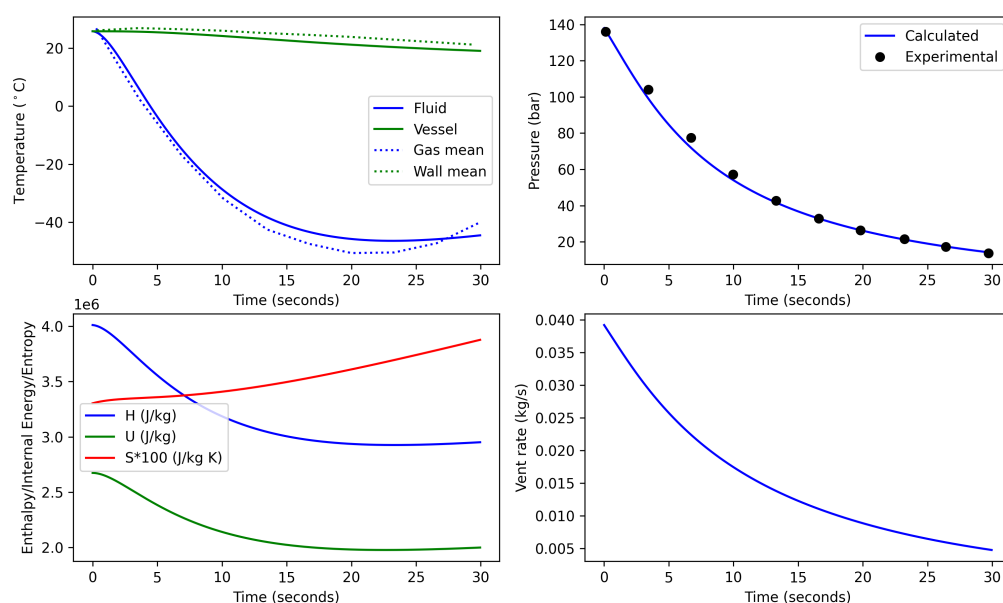


Figure 4.1: Calculations of hydrogen discharge experiment run 7 from (Byrnes, Reid, and Ruccia 1964). The figure shows calculated gas and wall temperature (full lines) compared to experiments (upper left), calculated and experimental pressure (upper right), specific thermodynamic state variables (lower left), and the calculated vent rate (lower right).

The figure layout is general for all the validation studies conducted and will be described in brief for eased readability. Each experiment figure is divided in 4 subplots; the upper left shows the simulated bulk gas temperature and vessel wall (average) temperature with measured values included with stipulated lines (if available), the upper right shows the simulated and measured pressure (if available, shown with points), the lower left shows the specific enthalpy, specific internal energy and the specific entropy of the bulk gas content as a function of time, the lower right shows the simulated mass rate leaving or entering the vessel.

As seen from Fig. 4.1, Fig. 4.2 and Fig. 4.3, the calculations generally compare well with the experimental results. It appears as though the faster depressurisation times are predicted the best, as seen from run 7 and 9. The much slower run 8 does not have as good a fit for the bulk gas temperature. However, on an absolute scale the results are still within 5°C of the experimentally determined gas temperature. The trend in *goodness of fit* could indicate that the heat transfer from the outside of the vessel to the gas lacks some details, which become detectable at the very slow depressurisation. For the fast depressurisation the heat transfer from the outside is of less importance and the inside heat transfer coefficient may be predicted better.

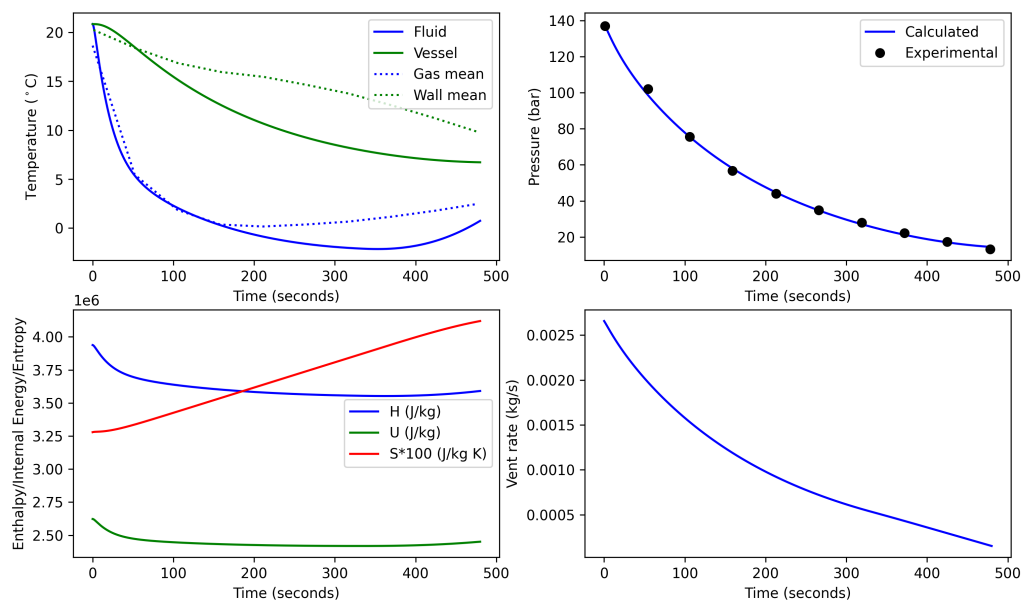


Figure 4.2: Simulation of run 8 from (Byrnes, Reid, and Ruccia 1964).

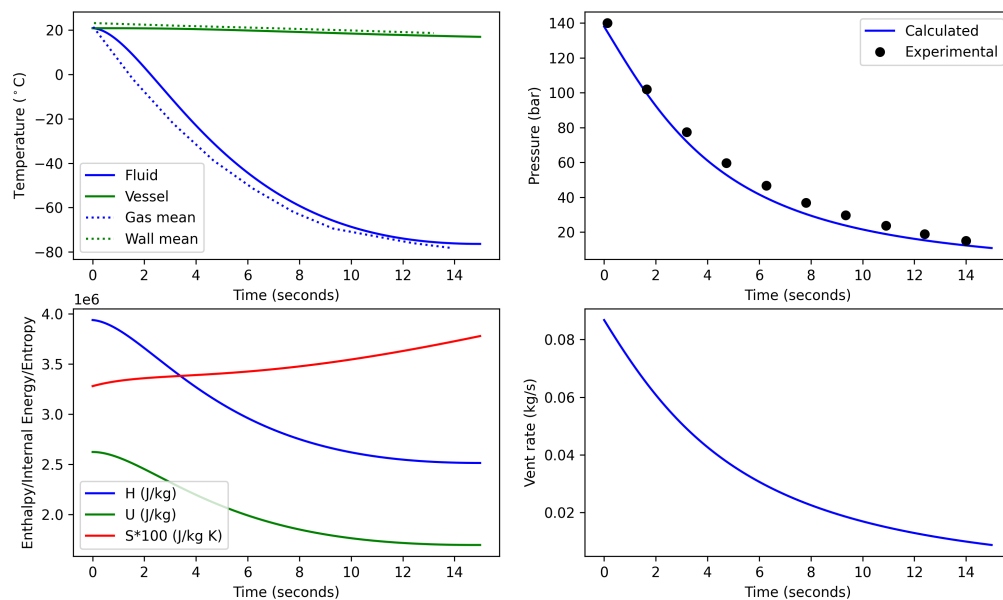


Figure 4.3: Simulation of run 9 from (Byrnes, Reid, and Ruccia 1964).

4.2 Hydrogen filling

In order to compare HydDown to experimental values of hydrogen filling operation the experiments reported by Striednig *et al.* (Striednig et al. 2014) are used. Experimental results were obtained using a type I tank (steel) with a volume of 0.0235 m³. The filling of hydrogen of gas into the vessel was carried out with an upstream reservoir kept at 350 bar and the flow was controlled with an electronically controlled dispenser. Vessel details are provided below.

Table 4.1: Key hydrogen vessel data (Striednig et al. 2014)

Parameter	Value
Mass	69..5 kg
Length	0.61 m
OD	0.280 m
Wall thickness	0.0129 m
Density	7740 kg/m ³
Heat capacity	470 J/(kg K)

(Striednig et al. 2014) reports many different experiments. In this comparison we will use experiments applying different pressurisation rates / durations using identical initial conditions. The experiments are reported in Fig. 6 and Fig. 7 in (Striednig et al. 2014).

In its current stage HydDown does not directly support specifying a fixed pressurisation rate expressed as a constant rate of change in pressure / pressure ramp rate (PRR). In order to emulate the experiments, an orifice model is used with an artificial high reservoir pressure, in order to ensure that the flow is critical during the entire pressurisation. This gives a constant mass rate, and translates to a fairly constant pressure ramp rate, although with real gas effects being clearly observable.

The results of the simulations with HydDown and the experiments from (Striednig et al. 2014) are shown in Fig. 4.4, Fig. 4.5, Fig. 4.6. As seen from the figures the experimental gas temperature appears to be simulated very accurately. The simulation with the fastest pressurisation under-predicts the experimental temperature slightly. The main features of the pressurisation temperature profile is a rapid increase in temperature, at the beginning of the pressurisation, where the effects of filling is dominating. At approx. 20-30 s the temperature increase levels off to a plateau, where the length of the plateau depends on the pressurisation rate. The slower the depressurisation, the longer the plateau. During this temperature plateau, the heat transfer from the gas to the vessel/surroundings balance

the temperature increase that would occur if performed at adiabatic (isolated) conditions. When the gas flow has ceased, at arrival of the final pressure, the temperature starts to drop due to cooling of the gas by the colder vessel wall driven by convective heat transfer. The temperature stabilises at a temperature which is higher than the initial/ambient temperature due to the heating of the steel vessel. The cooling of the vessel by ambient air is slower and would require a longer run time to be clearly visible.

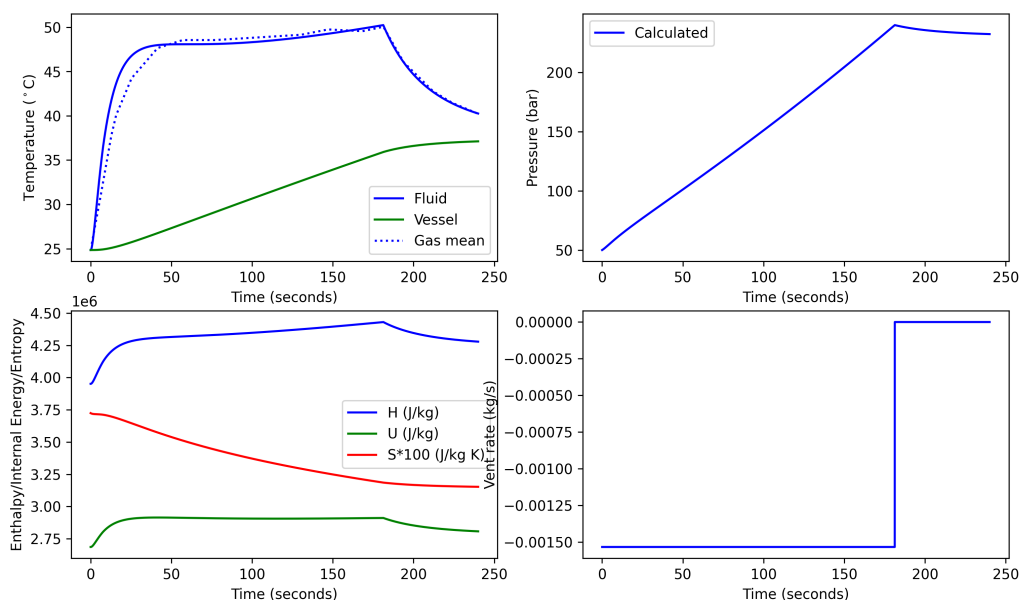


Figure 4.4: Simulation of H_2 pressurisation using 5 MPa/min (Striednig et al. 2014).

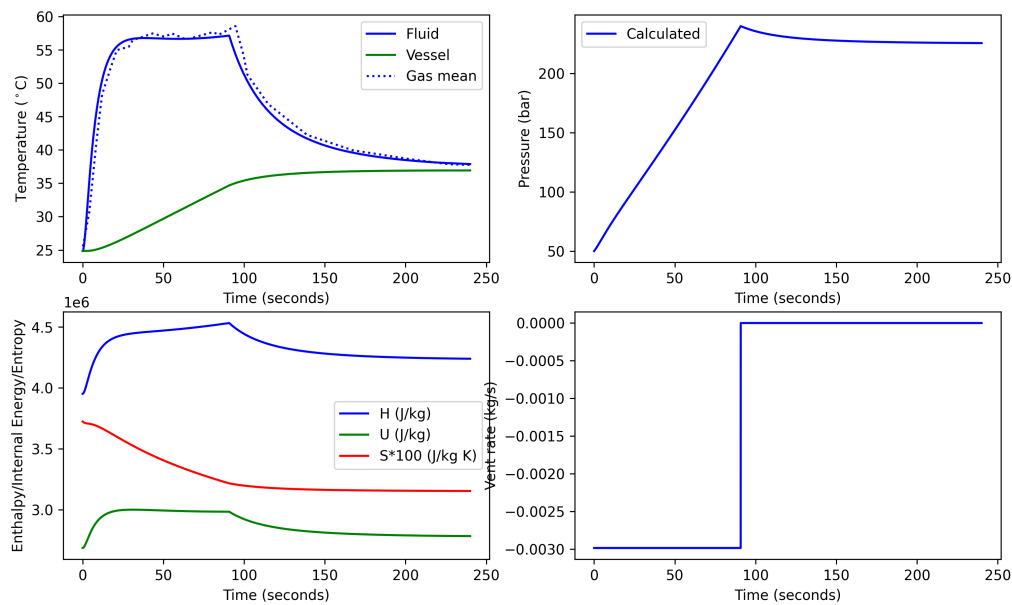


Figure 4.5: Simulation of H_2 pressurisation using 10 MPa/min (Striednig et al. 2014).

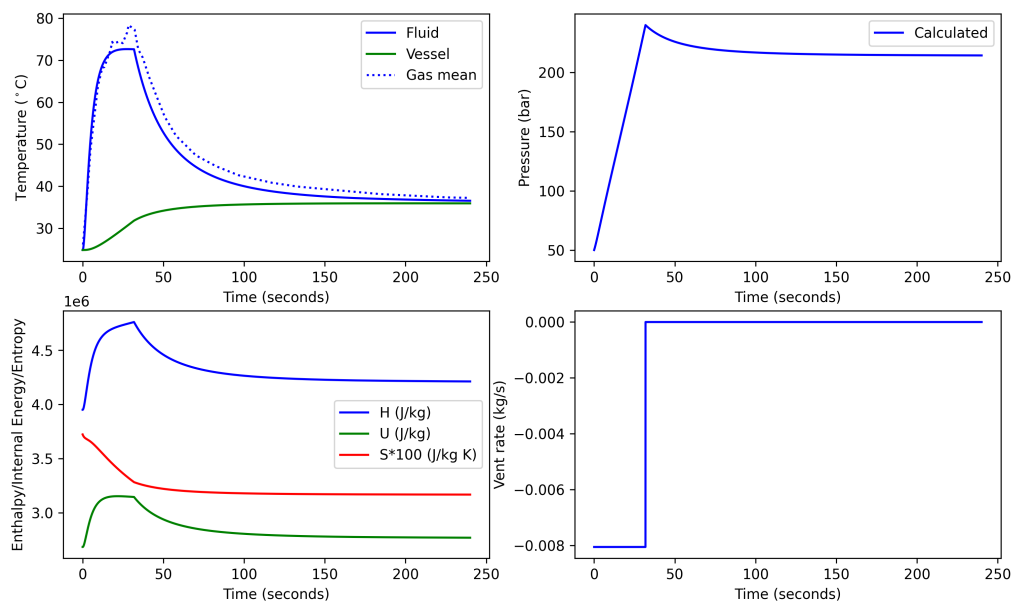


Figure 4.6: Simulation of H_2 pressurisation using 30 MPa/min (Striednig et al. 2014).

4.3 Nitrogen discharge

Calculations with HydDown is compared to experiment I1 from ref. (Haque et al. 1992). The experiment is a blowdown of a vertically oriented cylindrical vessel with flat ends. The vessel length is 1.524 m, the inside diameter is 0.273 m and the wall thickness is 25 mm. The vessel is filled with N_2 at 150 bar, at 15°C. Ambient temperature is 15°C. The blowdown orifice diameter is 6.35 mm. The results are shown in Fig. 4.7. The discharge coefficient of the orifice has been set to 0.8 in order to match the vessel pressure profile. The back pressure is set to atmospheric conditions.

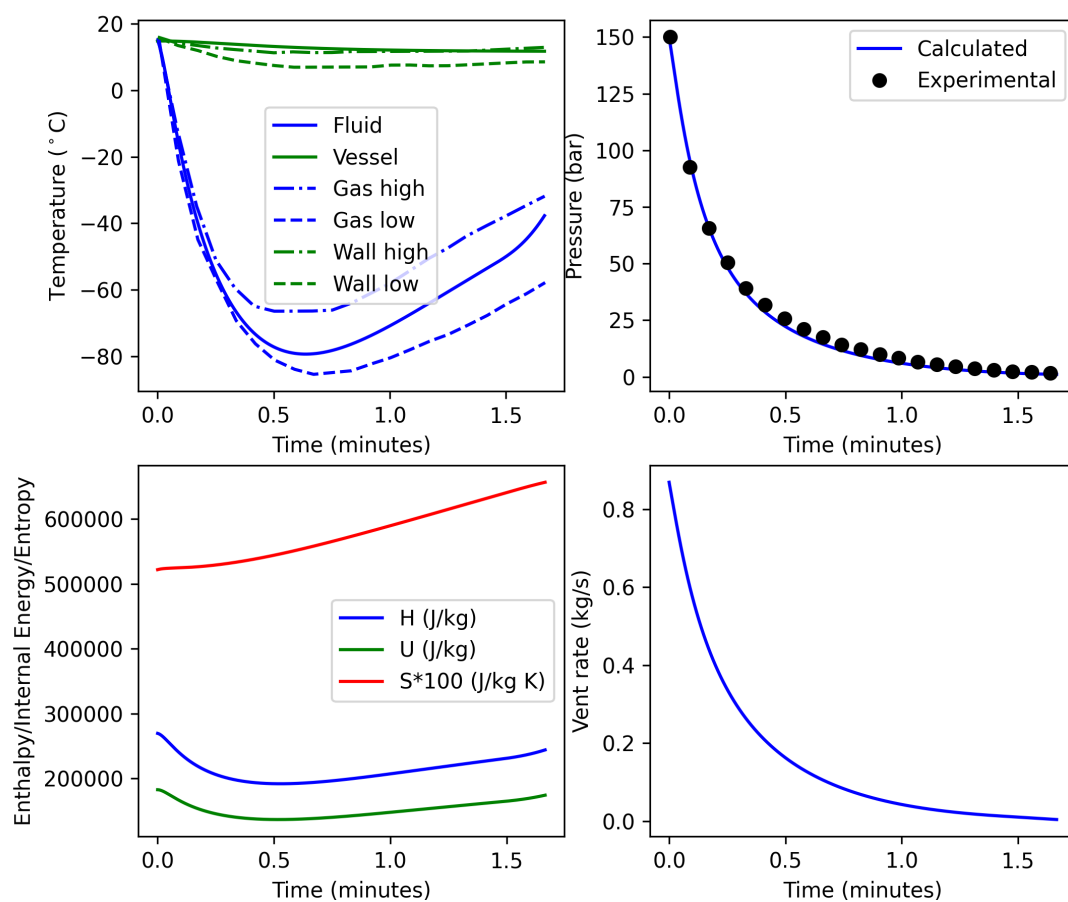


Figure 4.7: Calculations of nitrogen discharge emulating experiment I1 from (Haque et al. 1992). The figure shows calculated gas and wall temperature (full lines) compared to experiments (upper left), calculated and experimental pressure (upper right), specific thermodynamic state variables (lower left), and the calculated vent rate (lower right).

As seen from Fig. 4.7, the calculations compare well with the experimental results. The calculated temperature of the bulk vapor is within the experimental range of measured temperature at all times during the simulation. It is also noted that the minimum temperature is reached at approx. the same time as in the experiments. The calculated vessel inner wall temperature does not decline as rapidly as the experiments—but from around a calculation time of 60 s, the temperature is within the experimentally observed inner wall temperature. The main reason for the inability to match the vessel wall temperature is that the model ignores the temperature gradient from the outer to the inner wall surface and uses an average material temperature. Especially at the beginning of the discharge it is considered likely that a significant temperature gradient will exist.

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